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Original Article

Li₂AGeO₄ (A = Zn, Mg): Two novel low-permittivity microwave dielectric ceramics with olivine structureChunchun Li^{a,b,*}, Huaicheng Xiang^a, Minyu Xu^a, Ying Tang^a, Liang Fang^{a,*}^a State Key Laboratory Breeding Base of Nonferrous Metals and Specific Materials Processing, Guangxi Universities Key Laboratory of Non-Ferrous Metal Oxide Electronic Functional Materials and Devices, College of Material Science and Engineering, Guilin University of Technology, Guilin, 541004, China^b College of Information Science and Engineering, Guilin University of Technology, Guilin, 541004, China

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ABSTRACT

Two low-permittivity dielectric materials Li₂AGeO₄ (A = Zn, Mg) were prepared via the solid-state reaction method. X-ray diffraction analysis and Rietveld refinement indicated that both ceramics crystallize in an orthorhombic olivine structure with a space group *Pmn2₁*. Dense ceramics with high relative density and homogeneous microstructure were obtained. Li₂ZnGeO₄ densified at 1200 °C possessed a relative permittivity $\epsilon_r = 6.5$, a quality factor $Q \times f = 35,400$ GHz, and a temperature coefficient of resonant frequency. Li₂MgGeO₄ exhibited $\epsilon_r = 6.1$, $Q \times f = 28,500$ GHz, and $\tau_f = -74.7$ ppm/°C when sintered at 1220 °C. Additionally, the large negative τ_f values of Li₂AGeO₄ (A = Zn, Mg) ceramics were successfully adjusted compensated by forming composite ceramics with CaTiO₃ and near-zero τ_f values of +2.9 ppm/°C and +5.8 ppm/°C were achieved in 0.92Li₂ZnGeO₄-0.08CaTiO₃ and 0.90Li₂MgGeO₄-0.10CaTiO₃, respectively.

1. Introduction

The rapid advance in wireless communications to millimeter waves such as electronic transport cashing (ETC), ultra-high speed LAN, car anti-collision system on the intelligent transport system (ITS), in recent years, has triggered tremendous development in microwave dielectrics. In millimeter-wave applications, a large number of information can be transferred at a high speed [1–4]. Thus, three primary figures of merit are required: low permittivity (ϵ_r), high quality factor (low dielectric loss, $Q = 1/\tan\delta$), and near-zero temperature coefficient of resonant frequency (τ_f) [5,6]. Importantly, the signal delay time is proportional to the square root of permittivity according to the formula $t_d = \sqrt{\epsilon_r} l_e / c$ (where t_d , l_e , and c represents signal delay time, transmission distance, and velocity of light) [7]. Thus, low ϵ_r is desirable for high-speed transmission. Additionally, low permittivity also minimizes the cross-coupling between the substrates with the conductors and decreases the latency in long-range reception. Therefore, the development of low-permittivity microwave dielectric materials to cater high-speed transmission demands is becoming a hot issue.

In the past several decades, numerous low- ϵ_r materials, such as Al₂O₃, spinel MAl₂O₄ and M₂TiO₄ (M = Mg, Zn), and forsterite M₂SiO₄ (M = Mg, Zn), have been developed for millimeter-wave applications [8–14]. For example, Mg₂SiO₄ ceramics sintered at 1550 °C possessed excellent microwave dielectric performance with $\epsilon_r = 7.50$,

$Q \times f = 114,730$ GHz, and $\tau_f = -59$ ppm/°C [13]. The olivine structure has a general formula A₂BO₄, in which B⁴⁺ locates at the center of isolated tetrahedral geometry. Compared with spinel structure AB₂O₄, where A ions, usually divalent, are coordinated by four oxygen ions forming AO₄ tetrahedral, while B ions, usually trivalent, are coordinated by six oxygen ions forming BO₆ octahedral. Substitution of two trivalent B ions in the spinel AB₂O₄ by one divalent A and one tetravalent ion will result in an inverse spinel structure, such as Mg[MgTi]O₄ [15,16,17]. More recently, considerable attention has been focused on the M₂GeO₄ (M = Mg, Zn) olivine compounds due to their low permittivity and high quality factor [18–20]. For example, orthorhombic Mg₂GeO₄ ceramic sintered at 1450 °C had a low permittivity $\epsilon_r \sim 5.48$, a quality factor $Q \times f \sim 11,037$ GHz, and a temperature coefficient of resonant frequency $\tau_f \sim -27.61$ ppm/°C [19]. Nevertheless, these low permittivity microwave dielectric ceramics generally have high sintering temperatures and negative large τ_f values, which are undesirable for practical applications. In recent years, Li-containing oxides have been reported as low-firing ceramics, e.g. Li₄Ti₅O₁₂ [21], Li₂ZnTi₃O₈ [22], and Li₂ZnGe₃O₈ [23]. Therefore, it is possible to reduce the sintering temperature of M₂GeO₄ by introducing Li⁺ ionic into the olivine structure. Therefore, in the parent M₂GeO₄, we introduced 2 Li⁺ ions to substitute 1 M cation to synthesize two olivine Li₂MGeO₄ (M = Zn, Mg) ceramics with structure. And their microwave dielectric properties were characterized for the first time. In

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addition, CaTiO_3 ($\tau_f \sim +800$ ppm/ $^\circ\text{C}$) was utilized to form composite ceramics with $\text{Li}_2\text{MgGeO}_4$ to obtain thermally stable dielectric materials.

2. Experimental procedure

The Li_2AGeO_4 ($A = \text{Zn, Mg}$) ceramics were prepared by solid-state reaction of Li_2CO_3 (99.99%, Guo-Yao Co. Ltd., China), ZnO (99.99%, Guo-Yao Co. Ltd., China), MgO (99.99%, Guo-Yao Co. Ltd., China), and GeO_2 (99.999%, Guo-Yao Co. Ltd., China) as starting materials. Importantly, MgO powders were heated at 900°C for 1 h to remove moisture. The starting materials were weighed according to the stoichiometry and mixed through ball milling using ZrO_2 balls with ethanol as the medium for 6 h. The dried mixtures were calcined at 1100°C for 6 h in a platinum crucible, and re-milled for 4 h. Then PVA was added as binder, ground and pressed into 10 mm-diameter and 6 mm-height disks under a pressure of 350 MPa. The green compacts were firstly fired at 550°C in air for 2 h to expel the organic binder and then sintered at 1140°C – 1240°C for 6 h.

The crystal structure and phase purity were examined by powder X-ray diffraction (XRD) which was performed on a Panalytical X'Pert PRO diffractometer with $\text{Cu K}\alpha$ radiation. The surface microstructures of the pellets were observed by a Hitachi (Tokyo, Japan) S4800 scanning electron microscopy (SEM). The densities of the sintered ceramics were measured using the Archimedes' method. The permittivity and quality factor values were measured in the TE_{011} mode according to the Hakki-Coleman method using an Agilent N5230A network analyzer (Palo Alto, CA). The temperature coefficient of resonant frequency τ_f value was measured from 25°C to 85°C using a temperature chamber (Delta 9039, Delta Design, San Diego, CA).

3. Results and discussion

Fig. 1 shows the Rietveld refinements using Fullprof software were performed on the XRD data of $\text{Li}_2\text{ZnGeO}_4$ and $\text{Li}_2\text{MgGeO}_4$ and the calculated and measured diffraction patterns along with the unit cell volume and the schematic crystal structure. The consistent match between the calculated and measured diffraction patterns and the low residual factors (shown in the inset of Fig. 1(a)) confirm the phase purity of Li_2AGeO_4 . All the observed patterns that matched well with the standard JCPDS PDF No. 24-0675 for $\text{Li}_2\text{ZnGeO}_4$ and JCPDS No. 24-0627 for $\text{Li}_2\text{MgGeO}_4$ ceramics were successfully indexed with no secondary phase peaks detected, which indicates that both compounds crystallized in an orthorhombic olivine structure with a space group $\text{Pmn}2_1$. The refined unit cell volumes of the ceramics sintered at different temperatures are shown in Fig. 1(c). Obviously, with increasing sintering temperature, the unit cell volume decreased monotonously. This suggests that the distances between the ions were shortened and the crystal structure of Li_2AGeO_4 ($A = \text{Zn, Mg}$) were compressed with increasing sintering temperature. In addition, it is evidence that $\text{Li}_2\text{ZnGeO}_4$ (175.14 \AA^3) has a larger unit cell than $\text{Li}_2\text{MgGeO}_4$ (174.98 \AA^3), which can be explained by the larger ionic radius of Zn^{2+} (0.74 \AA) than Mg^{2+} (0.72 \AA). The theoretical densities were calculated based on the refined unit cell volume and atomic weight and yielded 4.08 g/cm^3 for $\text{Li}_2\text{ZnGeO}_4$ and 3.31 g/cm^3 for $\text{Li}_2\text{MgGeO}_4$. A schematic crystal structure of the Li_2AGeO_4 ($A = \text{Zn, Mg}$) is presented in Fig. 1(d). Li and Ge cations occupy the tetrahedral sites, with the neighboring $[\text{LiO}_4]$ and $[\text{GeO}_4]$ corner connected. The A cations (Zn or Mg) are coordinated by 6 oxygens to form $[\text{AO}_6]$ octahedra and all the octahedra are corner shared with the adjacent tetrahedra.

SEM images of the polished and thermally etched surfaces of the best-sintered samples are shown in the inset of Fig. 2. A dense and

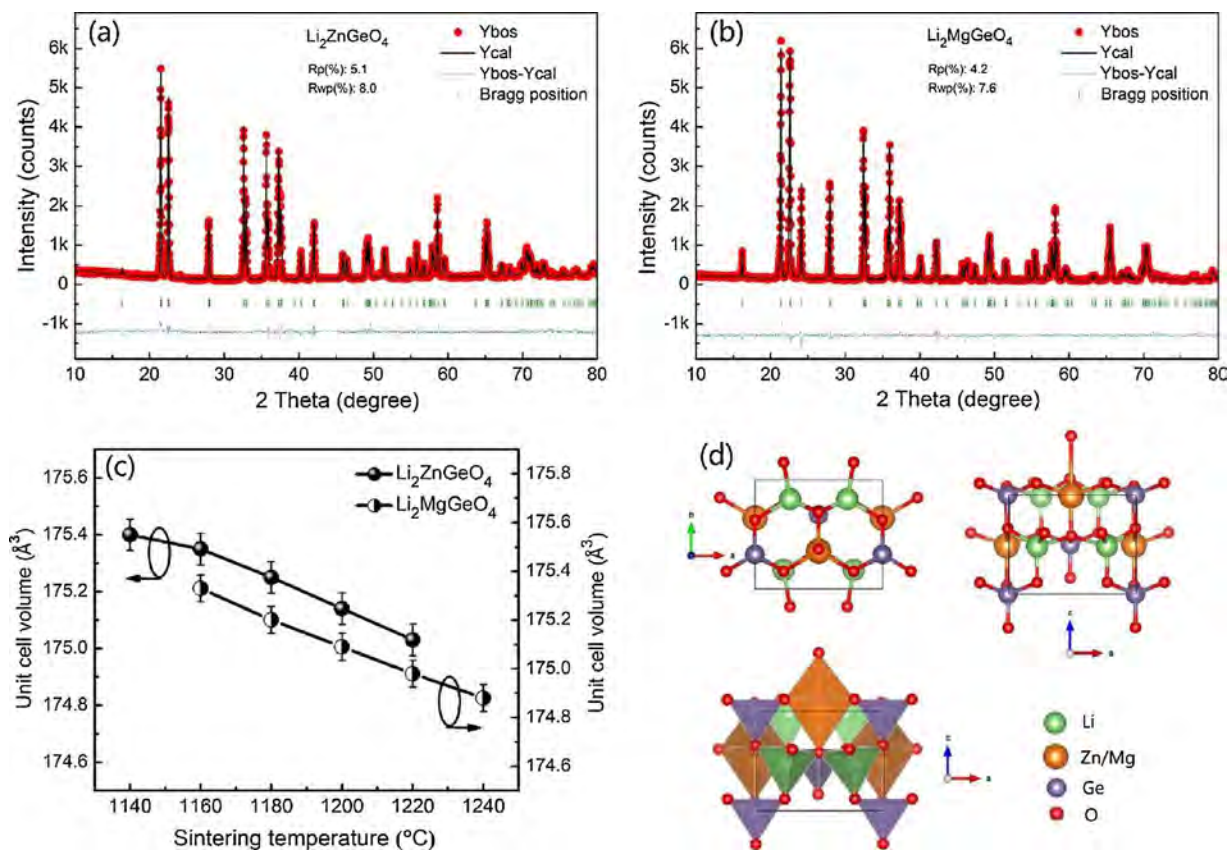


Fig. 1. Rietveld refinement of the room temperature XRD data for the (a) $\text{Li}_2\text{ZnGeO}_4$ and (b) $\text{Li}_2\text{MgGeO}_4$ ceramic, (c) unit cell volume of Li_2AGeO_4 ceramics as a function of temperature, and schematic crystal structure of the Li_2AGeO_4 ($A = \text{Zn, Mg}$) ceramics.

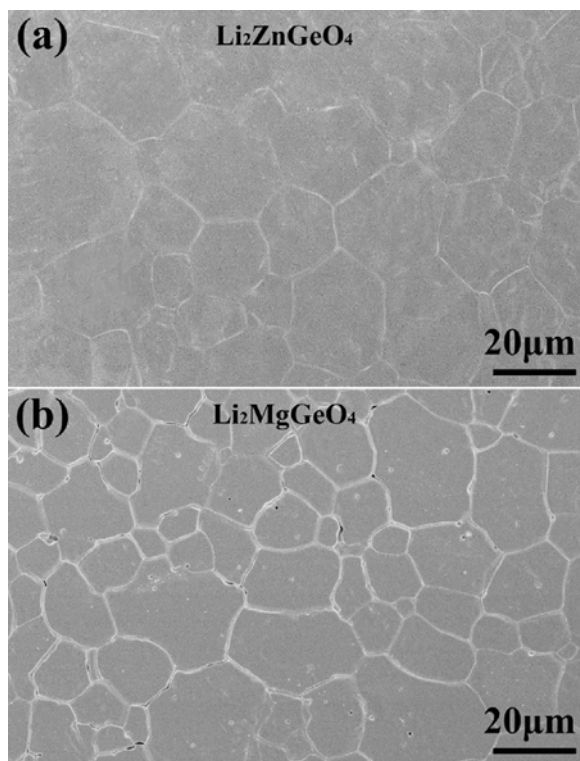


Fig. 2. SEM images of the polished and thermally etched surfaces of the Li_2AGeO_4 (A = Zn, Mg) ceramics sintered at optimum temperature.

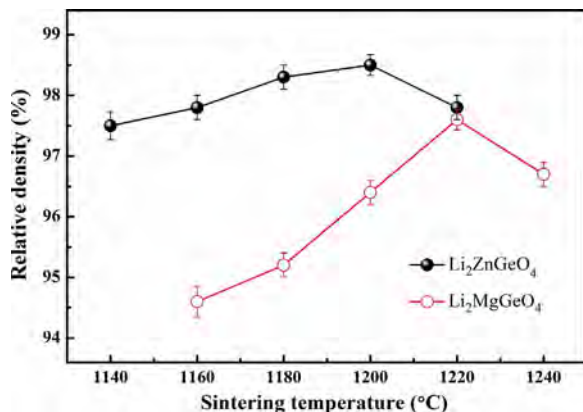


Fig. 3. The relative density of the Li_2AGeO_4 (A = Zn, Mg) ceramics as a function of temperature.

homogeneous microstructure with clear grain boundary is observed for both samples. The average grain size is 15 ~ 25 μm for $\text{Li}_2\text{ZnGeO}_4$ and 10 ~ 20 μm for $\text{Li}_2\text{MgGeO}_4$ ceramics, respectively. The variation in relative densities of the Li_2AGeO_4 (A = Zn, Mg) ceramics as a function of temperature is shown in Fig. 3. As the sintering temperatures increased, the relative density increased and obtained a saturated value, followed by a slight decrease in higher temperatures. The highest value ~ 98.5% of the theoretical density (~4.08 g/cm³) was obtained for $\text{Li}_2\text{ZnGeO}_4$ ceramic sintered at 1200 °C and ~ 97.6% for $\text{Li}_2\text{MgGeO}_4$ ceramic at 1220 °C. By comparison, the densification sintering temperature of the $\text{Li}_2\text{ZnGeO}_4$ is slightly lower than $\text{Li}_2\text{MgGeO}_4$ ceramic, but the relative density is higher than $\text{Li}_2\text{MgGeO}_4$ ceramic.

Microwave dielectric properties (ϵ_r , $Q \times f$, and τ_f) of the Li_2AGeO_4

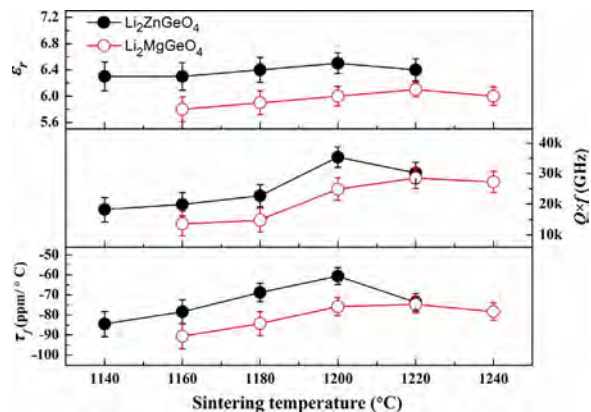


Fig. 4. The microwave dielectric permittivity, $Q \times f$ and τ_f value of the Li_2AGeO_4 (A = Zn, Mg) ceramics as a function of temperature.

(A = Zn, Mg) ceramics sintered at different temperatures are presented in Fig. 4. Over the sintering temperature range 1140–1220 °C, the relative permittivity of $\text{Li}_2\text{ZnGeO}_4$ ceramics varied slightly and fluctuated around 6.5. The ϵ_r of $\text{Li}_2\text{MgGeO}_4$ ceramics exhibited similar variation tendency to that of $\text{Li}_2\text{ZnGeO}_4$ ceramics. A maximum ϵ_r value 6.1 was achieved when sintered at 1220 °C for $\text{Li}_2\text{MgGeO}_4$, which was slightly lower than $\text{Li}_2\text{ZnGeO}_4$. The difference in relative permittivity could be explained by the lower ionic polarization of $\alpha(\text{Mg}^{2+})$ (1.32 Å³) than $\alpha(\text{Zn}^{2+})$ (2.04 Å³) based on the Clausius-Mosotti equation [24]. The calculated theoretical permittivity of $\text{Li}_2\text{ZnGeO}_4$ is 7.2 and 6.1 for $\text{Li}_2\text{MgGeO}_4$.

It is well known that at microwave frequency region the dielectric loss could be influenced by many factors that can be classified into two categories: the intrinsic factors related to the inharmonic lattice vibration and the extrinsic ones, e.g. density, grain boundary, pores, secondary phase, etc [25,26]. Obviously, $Q \times f$ values of Li_2AGeO_4 (A = Zn, Mg) ceramics exhibit similar variation trend to the corresponding relative densities. The result suggests that in the present work, the bulk density plays prominent roles in controlling the dielectric loss. Optimal $Q \times f$ values of 35,400 GHz and 28,500 GHz for $\text{Li}_2\text{ZnGeO}_4$ and $\text{Li}_2\text{MgGeO}_4$ ceramics were obtained at their relative densification temperatures (1200 °C and 1220 °C, respectively). Similarly, the temperature coefficient of the resonance frequency of Li_2AGeO_4 (A = Zn, Mg) ceramics of varied slightly and fluctuated with the sintering temperature. The τ_f values of -60.6 ppm/°C for $\text{Li}_2\text{MgGeO}_4$ and -74.7 ppm/°C for $\text{Li}_2\text{ZnGeO}_4$ were obtained. Unfortunately, the large negative τ_f values are unacceptable for practical applications where a near-zero τ_f value is expected. In order to adjust the thermal stability of Li_2AGeO_4 ceramics, an effective and successful strategy of an addition of another phase with opposite τ_f value to form composite ceramics is utilized [27,28]. In the present work, CaTiO_3 with a positive τ_f of +800 ppm/°C was chosen as the τ_f adjuster.

Fig. 5 shows the XRD patterns of $(1-x)\text{Li}_2\text{AGeO}_4 \cdot x\text{CaTiO}_3$ (A = Zn, Mg) ceramics. Only the peaks of CaTiO_3 (JCPDS No. 76-2400) and Li_2AGeO_4 phase were observed. This reveals that no chemical reaction occurred between Li_2AGeO_4 (A = Zn, Mg) and CaTiO_3 and they could co-exist in the $(1-x)\text{Li}_2\text{AGeO}_4 \cdot x\text{CaTiO}_3$ (A = Zn, Mg) system.

Fig. 6 shows the ϵ_r , $Q \times f$, and τ_f values of $(1-x)\text{Li}_2\text{AGeO}_4 \cdot x\text{CaTiO}_3$ (A = Zn, Mg) ceramics as a function of volume fraction of CaTiO_3 . As shown, the ϵ_r values of $(1-x)\text{Li}_2\text{AGeO}_4 \cdot x\text{CaTiO}_3$ (A = Zn, Mg) ceramics increased with increasing CaTiO_3 content, attributed to the higher ϵ_r value of CaTiO_3 (~170). However, CaTiO_3 addition to Li_2AGeO_4 (A = Zn, Mg) is detrimental to the $Q \times f$ value. As well known, the quality factor is sensitive to the extrinsic effects, especially the

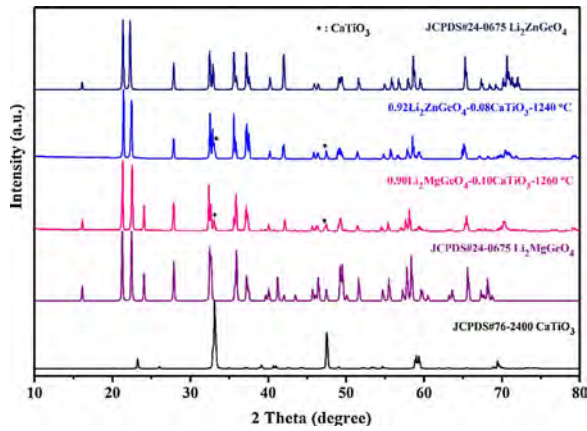


Fig. 5. XRD patterns of (a) 0.92Li₂ZnGeO₄-0.08CaTiO₃ ceramics sintered at 1240 °C, and (b) the 0.90Li₂MgGeO₄-0.10CaTiO₃ ceramics sintered at 1260 °C.

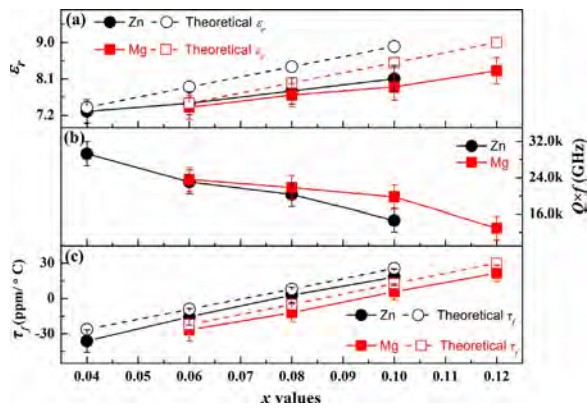


Fig. 6. The ϵ_r , theoretical ϵ_r , $Q \times f$, and theoretical τ_f values of $(1-x)\text{Li}_2\text{AGeO}_4-x\text{CaTiO}_3$ ($A = \text{Zn, Mg}$) ceramics as a function of the x values.

secondary phases. The CaTiO_3 addition as the second phase could deteriorate the $Q \times f$ values. Importantly, the τ_f values of $(1-x)\text{Li}_2\text{AGeO}_4-x\text{CaTiO}_3$ ceramics increased with the x values, showing a shift from negative to positive. As expected, a near-zero τ_f value of +2.9 ppm/°C was obtained in 0.92Li₂ZnGeO₄-0.08CaTiO₃ composite when sintered at 1240 °C and +5.8 ppm/°C for 0.90Li₂MgGeO₄-0.10CaTiO₃ composite sintered at 1260 °C. According to the empirical Lichtenecker logarithmic rule for a two-phase composite, the theoretical ϵ_r and τ_f values can be described by the following equations [29]:

$$\lg \epsilon = x_1 \lg \epsilon_1 + x_2 \lg \epsilon_2 \quad (1)$$

$$\tau_f = x_1 \tau_{f1} + x_2 \tau_{f2} \quad (2)$$

where x_1 and x_2 are the volume fractions, ϵ_1 and ϵ_2 are the permittivities, τ_{f1} and τ_{f2} are the τ_f values of the pure Li_2AGeO_4 ($A = \text{Zn, Mg}$) and CaTiO_3 phase, respectively. The theoretical ϵ_r and τ_f values of Li_2AGeO_4 - CaTiO_3 are calculated using the equations and shown in Fig. 6. By comparison, the measured ϵ_r and τ_f values were closed to the theoretical values.

Table 1 lists the sintering temperature and microwave dielectric properties of some low-permittivity microwave dielectric ceramics in the A_2BO_4 system. By comparison, the relative permittivity of the Li_2AGeO_4 ($A = \text{Zn, Mg}$) ceramics is close to that of the Zn_2MO_4 ($M = \text{Si, Ge}$). The quality factors of Li_2AGeO_4 are inferior compared with Zn_2MO_4 ($M = \text{Si, Ge}$), however, superior to $\text{Li}_2\text{MgSiO}_4$. Especially, it is worth noting that the Li-containing ceramics exhibited much lower sintering temperatures than others. Furthermore, all the listed A_2BO_4 microwave dielectric ceramics possess negative τ_f values. The successful thermal compensation in Li_2AGeO_4 ($A = \text{Zn, Mg}$) ceramics makes them possible candidates for practical applications.

4. Conclusions

Orthorhombic olivine structure Li_2AGeO_4 ($A = \text{Zn, Mg}$) ceramics were formed via the solid-state reaction method. The densification sintering temperature of the $\text{Li}_2\text{ZnGeO}_4$ ceramic is 1200 °C with relative density 98.5%, and at the same temperature to achieve the best microwave dielectric performances with $\epsilon_r \sim 6.5$, $Q \times f \sim 35,400$ GHz, and $\tau_f \sim -60.6$ ppm/°C. The $\text{Li}_2\text{MgGeO}_4$ ceramic sintered at 1220 °C for 6 h exhibited high relative density with 97.6% and excellent microwave dielectric properties with $\epsilon_r \sim 6.1$, $Q \times f \sim 28,500$ GHz, and $\tau_f \sim -74.7$ ppm/°C. Furthermore, the large negative τ_f values of Li_2AGeO_4 ($A = \text{Zn, Mg}$) ceramics could be compensated by forming composite ceramics with CaTiO_3 .

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Table 1
Microwave dielectric properties of the olivine structure ceramics.

Ceramics	S.T. (°C)	ϵ_r	$Q \times f$ (GHz)	τ_f (ppm/°C)	Reference
Mg_2SiO_4	1550	7.5	114,730	-59	[13]
Zn_2SiO_4	1325	6.6	198,400	-41.6	[14]
Mg_2GeO_4	1450	5.48	11,037	-27.61	[19]
Zn_2GeO_4	1300	6.87	102,700	-32.4	[20]
Mg_2SnO_4	1600	8.41	55,100	-62	[30]
$(\text{Mg}_{0.95}\text{Ni}_{0.05})_2\text{SnO}_4$	1550	8.72	103,100	-62.8	[31]
$\text{Li}_2\text{MgSiO}_4$	1250	5.1	15,380	-	[32]
$\text{Li}_2\text{MgGeO}_4$	1220	6.1 (± 0.3)	28,500 (± 200)	-74.7 (± 4)	This work
$\text{Li}_2\text{ZnGeO}_4$	1200	6.5 (± 0.3)	35,400 (± 300)	-60.6 (± 3)	This work
0.92 LZG-0.08 CT	1240	7.8 (± 0.4)	20,300 (± 200)	+2.9 (± 0.1)	This work
0.90 LMG-0.10 CT	1260	7.9 (± 0.4)	19,760 (± 200)	+5.8 (± 0.3)	This work

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